Polar auxin transport: models and mechanisms

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Summary
Spatial patterns of the hormone auxin are important drivers of plant development. The observed feedback between the active, directed transport that generates auxin patterns and the auxin distribution that influences transport orientation has rendered this a popular subject for modelling studies. Here we propose a new mathematical framework for the analysis of polar auxin transport and present a detailed mathematical analysis of published models. We show that most models allow for self-organised patterning for similar biological assumptions, and find that the pattern generated is typically unidirectional, unless additional assumptions or mechanisms are incorporated. Our analysis thus suggests that current models cannot explain the bidirectional fountain-type patterns found in plant meristems in a fully self-organised manner, and we discuss future research directions to address the gaps in our understanding of auxin transport mechanisms.

Key words: Computer simulation, Mathematical biology, Plant hormone

Introduction
Polar auxin transport
The plant hormone auxin plays a crucial role in the spatiotemporal control of plant development, and its patterns of distribution and activity must be tightly regulated. For example, in the root meristem a gradient of auxin with its maximum in the root tip determines the location of the quiescent centre and surrounding stem cells (Sabatini et al., 1999) and the regions where cell division, expansion and differentiation occur. Local maxima of auxin in the shoot apical meristem and in the differentiation zone of the mature root guide primordium outgrowth (Casimiro et al., 2001; Reinhardt et al., 2000), while in leaves, streams of auxin precede vein formation (Scarpella et al., 2006).

Much of the spatial distribution of auxin is caused by directional transport [polar auxin transport (PAT); see Glossary, Box 1]. The low pH in cell walls causes auxin to become protonated, allowing it to enter cells relatively easily. In addition, influx carriers of the AUX/LAX family pump auxin into cells. Owing to the higher pH in the cytosol, cellular auxin loses its ability to cross the membrane. Thus, auxin needs to be actively pumped out of cells by efflux carriers. Proteins of the PIN-FORMED (PIN) family are an important group of efflux carriers (Gälweiler et al., 1998; Muller et al., 1998, Paponov et al., 2005; Friml, 2003). They typically have a polar cellular distribution, leading to directed auxin transport across only those membranes where PINs are localised (hereafter, we describe PINs as ‘pointing’ in the direction along which auxin moves).
proteins constantly cycle to and from the PM, it is likely that this distribution of PIN proteins and thus implying a feedback (see Reinhardt et al., 2003), indicating its ability to alter the polar auxin can induce new primordia (Reinhardt et al., 2000; de Reuille et al., 2006; Wisniewska et al., 2006). In leaves, cell files with polar PIN distributions transport auxin along veins (Fig. 1A), whereas in the epidermis of the shoot apical meristem PINs point toward each auxin maximum (Reinhardt et al., 2003). Within these individual shoot primordia, PINs in the outer tissue layers are localised toward the auxin maximum in the primordium tip, whereas inner layers form veins with PINs pointing away from the maximum, connecting the primordium to the vasculature (Fig. 1C) – a pattern known as a reverse-fountain. Finally, in the root tip and lateral root primordia, an opposite PIN pattern (whereby PINs in the outer layers transport auxin away from the local maximum, whereas inner files of cells direct auxin towards the maximum) is observed (Fig. 1D), and this is referred to as a fountain (Blilou et al., 2005). Mathematical models for the root (Grieneisen et al., 2007) and shoot (de Reuille et al., 2006) meristem have demonstrated that the experimentally observed PIN polarity patterns are both necessary and sufficient for the correct build-up of auxin maxima.

PIN proteins undergo constant cycling to and from the plasma membrane (PM) (Geldner et al., 2001; Dhonukshe et al., 2007), allowing them to dynamically maintain their polarity and to quickly redistribute in response to endogenous triggers (primordia formation, gravitropic response) or external stimuli (wounding, stretching). It has been experimentally shown that externally applied auxin can induce new primordia (Reinhardt et al., 2000; Reinhardt et al., 2003), indicating its ability to alter the polar distribution of PIN proteins and thus implying a feedback (see Glossary, Box 1) loop between auxin and its own transport. As PIN proteins constantly cycle to and from the PM, it is likely that this feedback represents a regulatory effect of auxin on PIN cycling. Indeed, it has been shown (Paciorek et al., 2005; Robert et al., 2010) that ectopically added auxin counteracts the PIN internalisation induced by the exocytosis-inhibiting drug BFA, and this has been interpreted as an inhibitory effect of auxin on endocytosis. However, it is currently hard to establish whether intra- or extracellular auxin or auxin flux is affecting PIN cycling in these experiments, and whether cells sense auxin directly or also indirectly via mechanosensitive signalling pathways.

**PAT models**

Long before the discovery of PIN proteins and the auxin dependency of their polar localisation, it was already hypothesised that auxin positively influences its own transport and spatial distribution, implying that auxin patterning might be self-amplifying and potentially even self-organising (Sachs, 1969). It is this implication that served as a major inspiration for the numerous modelling studies in this area. In the case of self-amplification (see Glossary, Box 1), the plant is capable of responding to and enhancing a superimposed auxin prepattern, for example a local source or sink, but requires this prepattern to persist. By contrast, in the case of self-organisation (see Glossary, Box 1), a transient prepattern is sufficient for initialisation and subsequent autonomous maintenance of the formed pattern.

Owing to the current lack of a fully mechanistic molecular understanding of how cells sense auxin and how this subsequently influences PIN polarity and auxin transport, most current models have taken a top-down approach, correlating the observed auxin and PIN polarity patterns in the tissue of interest to derive a hypothetical feedback mechanism. Based on their proposed feedback mechanism, PAT models can be divided into two main classes: flux-based and concentration-based models.

Flux-based models are based on Sachs’ canalisation hypothesis (Sachs, 1969), which states that cells experiencing flux of auxin in a certain direction will increase their capacity to transport the molecule in that direction, and is based on the observation that, during vein formation, auxin transport channels become gradually more distinct. The auxin transport capacity is represented by membrane permeability in early models (Mitchison, 1980; Mitchison, 1981) and by membrane PIN concentration in later models (e.g. Fujita and Mochizuki, 2006; Feugier and Iwasa, 2006; Alim and Frey, 2010; Feugier et al., 2005; Stoma et al., 2008). Flux-based models have mainly been used to model venation patterns, and demonstrate that small fluctuations in auxin may be amplified into more distinct streams, with PINs pointing in the direction of the flux, i.e. with-the-flux (Fig. 1A, inset).

Concentration-based models (e.g. Smith et al., 2006; Newell et al., 2007; Jönsson et al., 2006; Merks et al., 2007) were formulated after the discovery of PIN proteins and therefore all explicitly model membrane PIN levels. In these models, PIN levels increase on the membrane facing the neighbouring cell with the highest auxin level, i.e. up-the-gradient. This proposed feedback mechanism was inspired by observations in the shoot apex, where PINs in the epidermal layer orient toward local auxin maxima that develop into organ primordia (Fig. 1B, inset) (Reinhardt et al., 2003). Concentration-based models are sufficient to obtain phyllotaxis-like patterns by amplifying small local increases in auxin into distinct maxima while simultaneously depleting neighbouring cells, resulting in the occurrence of new maxima at fixed distances from older maxima.

More recently, efforts have been made to construct PAT models capable of displaying both up-the-gradient phyllotaxis and with-
the-ﬂux venation types of PIN patterning (Bayer et al., 2009; Merks et al., 2007; Stoma et al., 2008) in order to explain reverse-fountain-type patterns (Fig. 1C). Additionally, somewhat more mechanistic feedback (see Glossary, Box 1) loops of auxin on PAT have been suggested. Based on the observation that PIN polarity correlates with stress-related microtubule alignment, a model has been proposed in which auxin inﬂuences wall stress, which in turn inﬂuences PIN localisation (Heisler et al., 2010). Alternatively, it has been proposed that PIN polarity is regulated by auxin receptors in the apoplast, which inhibit local PIN endocytosis after binding auxin (Wabnik et al., 2010).

**A framework to analyse and compare PAT models**

In this Hypothesis we will perform a detailed comparison and analysis of a broad range of models for auxin patterning in plant development. For ﬂexible developmental patterning, the feedback between active polar auxin transport, which generates auxin distribution patterns, and these auxin distributions in turn shaping the strength and direction of this polar auxin transport, is crucial. Therefore, we restrict this analysis to models that incorporate polar auxin transport and its auxin-dependent regulation, focusing on both the auxin and PIN distribution patterns that are generated. Thus, we will not include strictly mechanical models that do not consider feedback on auxin transport, Turing-type models (see Glossary, Box 1) in which only passive undirected auxin transport is considered, or models in which the directions of active PAT are assumed to be constant (e.g. Grieneisen et al., 2007).

One straightforward way to evaluate and compare the various PAT models is by analysing the tissue level auxin and PIN patterns that they generate and compare these with experimental data. A number of excellent reviews have been written that describe these efforts (Wabnik et al., 2011; Heisler and Jönsson, 2006; Kramer, 2009; Garnett et al., 2010). However, for a mathematical model to work, detailed speciﬁcations have to be made for the dynamics on the subcellular membrane segment, cellular and tissue levels.

Indeed, apart from the simple dichotomy in concentration and ﬂux-based models, a variety of mathematical formulations for processes such as PIN-mediated auxin pumping, PIN dynamics and auxin-PIN feedback functions are used in the different models, the precise relevance of which is not trivial for even the experienced modeller. Therefore, as recently pointed out in a review by Jönsson and Krupinski (Jönsson and Krupinski, 2010), it is of crucial importance to develop a general framework to analyse and compare these different mathematical formulations, classify them into a limited number of corresponding biological assumptions, and determine how these inﬂuence model patterning behaviour.

Another crucial aspect is to determine the extent to which the model generates these patterns in an autonomous self-organised manner. Although self-organising and self-amplifying models may produce similar spatial patterns, whether this patterning requires a persistent prepattern is of key relevance. If auxin patterning required a prepattern, then a mechanism would be needed to explain the generation of these prepatterns. Furthermore, changes in auxin patterns precede major changes in the expression of developmental genes and cell morphology. Thus, for plants to keep generating new organs robustly throughout their life, it seems essential that the new auxin patterns required by each newly forming organ primordium arise in a largely self-organised manner independently of a prepattern.

In this Hypothesis we develop a general mathematical and simulation framework in which we analyse and compare most currently published PAT models. We translate the various mathematical formulations that are used into biological assumptions, and analyse how these relate to model behaviour, investigating both the type of auxin patterns generated and to what extent these are fully self-organised. We restrict our analysis to assumptions regarding PIN-mediated auxin transport dynamics, PIN cycling and auxin-PIN polarisation feedback, assuming these to be the main determinants for auxin patterning. The transcriptional effect of auxin on PIN expression levels (Vieten et al., 2005) is
mostly ignored because the short-term effect of auxin on PIN polarisation is expected to be more important for the self-organising capabilities of a mechanism than the long-term effect of auxin on overall PIN levels. Model behaviour will be analysed with regard to three levels of increasing complexity: the membrane segment, single-cell and one-dimensional (1D) tissue levels (Fig. 2A-C, respectively). Our 1D tissue level analysis does not allow us to directly determine the type of 2D or 3D patterns that a model can generate. However, it allows us to determine whether PINs orient away or toward an auxin maximum and whether a single mechanism can generate the two opposite polarisation patterns observed in fountain-type patterns. Additionally, we can establish a model’s self-organising potential in simple 1D tissue simulations. If a model generates self-organised patterns in one dimension, it will also do so in two or three dimensions. Thus, we can still extrapolate our 1D results to qualitatively predict the extent to which a model can generate complex tissue patterns, such as fountain-like patterns, and to what extent this patterning might be self-organised.

Our analysis will demonstrate that most published models are capable of producing robust, self-organised auxin and PIN patterns, and that flux-based and concentration-based models achieve this with largely similar biological assumptions. It appears that flux-based models are somewhat more versatile in explaining different PIN orientations, but both model categories seem to have difficulties in explaining the bidirectional fountain-type patterns observed in planta in a fully self-organised manner. We thus conclude that none of the currently available models robustly produces fountain-like patterns by a single mechanism. Finally, future directions for research are recommended in order to close in on the mechanistic basis of auxin feedback in plants and determine how fountain-type patterns are generated.

**Model behaviour on three levels of organisation**

In the following sections, published PAT models will be compared in terms of the mathematical formulations used for PIN-mediated auxin pumping dynamics, PIN cycling and auxin-PIN feedback, the biological assumptions to which these correspond, and the resulting model behaviour on the membrane segment, single-cell and tissue levels. A major goal of our analysis is to determine the maximal self-organising potential (see Glossary, Box 1) of models and the mathematical and biological assumptions on which this potential critically depends. As model behaviour depends both on mathematical formulations and parameter settings, we will classify a model as self-organising if self-organising patterns occur in at least a region of parameter space. We will first analyse in detail two representative examples of PAT models (chosen solely to facilitate discussion of our analysis of self-organising properties), following which we apply the same evaluation to a broader range of published models.

We have developed a generalised mathematical framework to study model assumptions and resulting behaviour (Fig. 2). In our analysis, the original mathematical formulation and biological assumptions of the discussed models are followed, but some simplifications are used to allow for an analytical approach. For ease of comparison we introduce a single set of variable and parameter names used for all models, rather than adopting the different names used in the various publications (see Table 1 for these parameters and variables, their biological meanings and default values). An in-depth description of the mathematical framework is provided in supplementary material Appendix S1.

We start with a simplified model at the membrane segment level and extend this first to the single-cell and subsequently to the 1D tissue level. Boxes 2 and 3 explain the mathematical formulation and analysis of the membrane segment and single-cell models. With regard to membrane segments, we are interested in whether the model allows for two alternative stable equilibria (see Glossary, Box 1) – termed membrane bistability (see Glossary, Box 1) – in which case membrane segments have either a discrete high or low PIN level. Alternatively, membrane segments may have a single equilibrium PIN level, in which case the PIN levels on a membrane segment may vary only according to locally experienced auxin levels that influence the precise location of this equilibrium (‘graded distribution’ (Fig. 2A)).

At the cell level we are interested in whether the model allows for stable equilibria in which a cell has a distinct high PIN level on one membrane segment and a distinct low PIN level on the

<table>
<thead>
<tr>
<th>Variable</th>
<th>Biological meaning</th>
<th>Units</th>
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<tbody>
<tr>
<td>( P )</td>
<td>PIN concentration on membrane segment of interest</td>
<td>[]</td>
</tr>
<tr>
<td>( A )</td>
<td>Auxin concentration in neighbouring cell of interest</td>
<td>[]</td>
</tr>
<tr>
<td>( F )</td>
<td>Auxin flux over membrane segment of interest into neighbouring cell of interest</td>
<td>[ls(^{-1})]</td>
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<table>
<thead>
<tr>
<th>Parameter</th>
<th>Biological meaning</th>
<th>Units and default values</th>
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<tbody>
<tr>
<td>( p )</td>
<td>Auxin production</td>
<td>1 [ls(^{-1})]</td>
</tr>
<tr>
<td>( d )</td>
<td>Auxin decay</td>
<td>0.5 [s(^{-1})]</td>
</tr>
<tr>
<td>( \bar{p}_{\text{pax}} )</td>
<td>Passive influx over the membrane segment of interest into the neighbouring cell of interest</td>
<td>0.01 [ls(^{-1})]</td>
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<tr>
<td>( \bar{p}_{\text{pin}} )</td>
<td>PIN-mediated influx over the membrane segment of interest into the neighbouring cell of interest</td>
<td>1 [s(^{-1})]</td>
</tr>
<tr>
<td>( e )</td>
<td>Total efflux out of the neighbouring cell of interest over the membrane segment of interest</td>
<td>1 [s(^{-1})]</td>
</tr>
<tr>
<td>( \bar{e}_{\text{pax}} )</td>
<td>Passive efflux out of the neighbouring cell of interest over the membrane segment of interest</td>
<td>0.01 [s(^{-1})]</td>
</tr>
<tr>
<td>( \bar{e}_{\text{pin}} )</td>
<td>PIN-mediated efflux out of the neighbouring cell of interest over the membrane segment of interest</td>
<td>1 [s(^{-1})]</td>
</tr>
<tr>
<td>( k_{\text{on}} )</td>
<td>PIN exocytosis/recycling rate</td>
<td>–</td>
</tr>
<tr>
<td>( k_{\text{onp}} )</td>
<td>Basal exocytosis/recycling rate</td>
<td>0.01 [s(^{-1})]</td>
</tr>
<tr>
<td>( k_{\text{ont}} )</td>
<td>Auxin-dependent exocytosis/recycling rate</td>
<td>1 [s(^{-1})]</td>
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<tr>
<td>( k_{\text{out}} )</td>
<td>PIN endocytosis rate</td>
<td>1 [s(^{-1})]</td>
</tr>
<tr>
<td>( k_{\text{outp}} )</td>
<td>Basal endocytosis rate</td>
<td>–</td>
</tr>
<tr>
<td>( k_{\text{outf}} )</td>
<td>Auxin-dependent endocytosis rate</td>
<td>–</td>
</tr>
<tr>
<td>( p_{\text{tot}} )</td>
<td>The total amount of PINs in one cell</td>
<td>10 []</td>
</tr>
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Note that under default conditions feedback is on \( k_{\text{off}} \) with \( k_{\text{off}} \) constant. Consequently, \( k_{\text{off}} \) is not defined as a constant value, but is instead a function of \( k_{\text{onp}} \), \( k_{\text{ont}} \) and auxin level or flux, while \( k_{\text{off}} \) has a defined constant value and \( k_{\text{outp}} \) and \( k_{\text{outf}} \) are not defined. If, instead, feedback is on \( k_{\text{off}} \) the situation is reversed. [] is used for concentration.
At the membrane segment level, we take into account the model’s description of PIN (P) dynamics on a single membrane segment and auxin concentration (A) in the adjoining neighbouring cell (see Fig. 2). Other PIN and auxin concentrations are assumed to be constant, which constitutes a mathematical tool enabling our membrane level analysis by limiting the number of variables included in the model, rather than a biological assumption; indeed, as we move in our analysis to the single-cell and tissue level these concentrations are no longer assumed constant. As we do not explicitly model the cell wall, the efflux over the membrane segment equals the influx into the neighbouring cell and vice versa. Hence, we can write for the auxin concentration:

$$\frac{dA}{dt} = p + i_{in} + i_{ex} - eA - dA,$$

with parameter symbols, meanings and default values described in Table 2.

Eq. 2 describes the auxin equilibrium line, which is the series of points (A,P) for which the auxin concentration is in equilibrium (dA/dt=0):

$$P = \frac{(e+d)A - i_{ex}}{i_{in}},$$

which is the solid line in panel A. To the left of the equilibrium line, net influx occurs and auxin increases (→), whereas to the right, auxin decreases due to net efflux (←).

Taking into consideration the dynamic cycling of PIN proteins between membrane and cytosol, we can model membrane segment PIN levels as:

$$\frac{dp}{dt} = k_{on} - k_{off}P.$$  \hspace{1cm} (3)

If $k_{on}$ and $k_{off}$ are constant (no feedback), the equilibrium line for PINs (dP/dt=0) is given by $P=k_{on}/k_{off}$ (dashed line in B). Above this line, PINs decrease (↓), whereas below this line they increase (↑). A positive feedback of auxin flux or concentration on PIN levels is incorporated by allowing the auxin flux or level to either increase $k_{on}$ or decrease $k_{off}$. This dependence of $k_{on}$ or $k_{off}$ on auxin can be modelled using either a linear feedback (e.g. $y=ax$), a superlinear feedback (e.g. $y=ax^2$), a (sub)linear saturating feedback [e.g. $y=\alpha x/(\beta+x)$] or a superlinear saturating feedback [e.g. $y=ax^2/(\beta+x^2)$], depending on the assumed effect of auxin on PIN cycling dynamics (with $a$ and $b$ being arbitrary scaling parameters and $n$ a measure of cooperativity). A hypothetical PIN equilibrium line resulting from a saturating superlinear feedback is given in C.

In D and E the auxin equilibrium line and auxin dynamics are drawn together with the two PIN equilibrium lines and PIN dynamics, allowing us to assess the overall behaviour of the membrane segment model. In the figures we can find the system’s equilibrium points (circles), in which both $dA/dt=0$ and $dP/dt=0$ and hence no changes occur, as intersection points of the A and P equilibrium lines. In addition, we can determine the stability of the equilibria from the direction of the dynamics (arrows) near an equilibrium. An equilibrium is stable if all arrows point toward it (black circles), whereas an equilibrium is unstable if one or more arrows point away from it (white circles). In D, the system has a single stable equilibrium. In E, the system has three equilibria, the middle one being unstable and the outer two stable, and thus represents a bistable membrane segment system.

Tissue level behaviour is analysed using a 1D array of five cells (Fig. 2C). First, the cells are organised as a file with a source of auxin at one end and a sink at the other. This allows us to test whether cells polarise and to determine the direction of polarity with respect to the global auxin gradient (i.e. the type of pattern formed) (Fig. 2C). Second, the file is wrapped into a ring, and a transient increase in the auxin level of a single cell is applied. In this system, we can investigate whether a transient perturbation spreads out and leads to persistent patterning or dies out, causing the tissue to return to its uniform state. If persistent pattern formation occurs in both configurations (file and ring) the tissue pattern is self-organising. If patterning is self-organised and driven by lower level membrane bistability and cell polarity, we will name it polarity-driven self-organisation (see Glossary, Box 1) to distinguish it from more conventional and well-analysed Turing-like self-organised patterning. If patterning is not self-organised and is instead maintained only in the presence of sources and sinks, the pattern formation is considered as gradient-driven amplification (see Glossary, Box 1), i.e. to be self-amplifying.

**Analysis of published models**

**Example 1: flux-based model**

Stoma et al. (Stoma et al., 2008) developed a flux-based PAT model in which the PIN concentration at a given membrane segment is assumed to increase with the net efflux of auxin over that membrane segment. Both linear and quadratic feedback functions for the dependence of PIN levels on auxin were used. Linear feedback (see Glossary, Box 1) implies that similar flux increases cause similar PIN level increases independently of the flux level. By contrast, superlinear (e.g. in this case quadratic) feedback (see Glossary, Box 1) implies that similar flux increases cause larger PIN increases for higher flux levels. The authors furthermore assumed that the availability of PIN proteins within a single cell is never limiting, i.e. the allocation of PINs to membrane segments does not lower the amount of PINs available for exocytosis/recycling sufficiently to cause competition between membrane segments. Finally, pumping of auxin by PIN proteins is assumed to depend linearly on intracellular auxin concentrations. In biological terms this assumption requires that PIN proteins are sufficiently available to handle large auxin concentrations and are thus not limiting for the rate of flux.
Development 140 (11)

Box 3. The single cell

To model a single cell in a simplified manner we consider the PIN levels on the two opposing membrane segments of the cell (P0 and P1) and the auxin concentrations in the two neighbouring cells (A0 and A1) (see Fig. 2). Again, we assume other auxin and PIN concentrations to be constant. The cell level model thus can be written as:

\[
\frac{dA_i}{dt} = p + j_{pas} + j_{pin}P_i - eA_i - dA_i, \quad (4a)
\]

\[
\frac{dP_i}{dt} = k_{on} - k_{off}P_i, \quad (4b)
\]

with \(i=0\) or \(1\).

If, in a model, a limiting PIN pool is assumed, the equation for \(P_i\) becomes:

\[
\frac{dP_i}{dt} = k_{on}(P_{tot} - P_i) - k_{off}P_i. \quad (5)
\]

The recycling rate now depends on the amount of available PINs in the cytosol, which is the total amount of PINs per cell (\(P_{tot}\)) minus the PINs that are bound to membrane segments.

To simplify the system and allow for 2D phase plane (see Glossary, Box 1) analysis, we make a quasi-steady-state assumption by letting the auxin dynamics at all times be in equilibrium with the amount of PINs at the membrane segment (\(dA_i/dt=0\)). This assumption does not alter the model behaviour in which we are interested, namely the number of equilibria. Hence, the expression for \(A_i\) becomes:

\[
A_i = \frac{P + j_{pas} + j_{pin}P_i}{e + d}. \quad (6)
\]

This leaves us with two equations for the PINs (Eq. 4b or Eq. 5) in which either \(k_{on}\) or \(k_{off}\) rates depend on auxin flux or concentration and hence on the equilibrium values of auxin given by Eq. 6.

The behaviour of the cell model depends both on the behaviour of the model at the membrane segment level and on whether a limiting PIN pool is assumed. If the PIN pool is non-limiting, \(P_0\) and \(P_1\) are completely independent of each other, and cell level equilibria arise from all possible combinations of membrane segment equilibria. In the case of a single membrane segment equilibrium, cell level equilibrium lines intersect once in a single stable equilibrium (panel A, solid line for \(P_0\), dashed line for \(P_1\)) and differences between \(P_0\) and \(P_1\) can only arise if the two membranes persistently experience different auxin concentrations or fluxes, thus simply shifting the location of this single equilibrium. By contrast, in the case of a bistable membrane segment, cell level equilibrium lines intersect nine times, producing a total of four stable equilibria: two alternative polar states, an apolar rest state, and a bipolar state (B). If, by contrast, \(P_0\) and \(P_1\) do influence each other via competition for a limiting PIN pool, different situations arise. In the case of a membrane segment with a single equilibrium there is most likely still only one, symmetrical, equilibrium at the single-cell level (A). In the case of a bistable membrane segment model there may be only two stable apolar equilibria (C), or an additional apolar rest state, in which \(P_0\) and \(P_1\) are equal and low (D), but no bipolar state is present.

The authors demonstrated that, in an otherwise homogeneous tissue, persistent sinks attract a small initial flux, causing PINs to orient toward them and auxin patterns to be built up. They showed that, if the degradation of auxin in sinks is sufficiently fast, auxin concentrations become lowest in the sinks, producing with-the-flux but up-the-gradient patterning as expected for with-the-flux auxin and PIN concentrations (Fig. 3B).

When we couple five cells into a cell file flanked by a source and a sink, we observe the expected with-the-flux PIN localisation in the direction of the sink, as in the original paper (Fig. 3C). Furthermore, the strength of the sources and sinks does not influence either the ability to polarise or the strength of polarisation (Fig. 3C, upper two cell files). Finally, for parameter settings in which auxin decay is very slow and hence auxin concentrations are higher in the sinks rather than in the source, we also observed up-the-gradient behaviour (Fig. 3C, bottom cell file). PIN polarity can also be formed and maintained in a ring of cells that only receives a transient perturbation in auxin levels (Fig. 3C). Note that as all cells in the ring pump auxin to their right and receive auxin from two separate regions of membrane segment behaviour, but not two stable equilibria. Owing to the absence of a limiting PIN pool, membrane PIN levels are fully independent of each other and only depend on local flux. Consequently, given a bistable membrane segment, a single cell can have four equilibria that are combinations of the two equilibria at each membrane segment. There is a single stable equilibrium in which both membrane segments have the same low PIN level, which we will refer to as the apolar rest state. In addition to this, there are three unstable equilibria, separating regions for which either one or both of the membrane segments obtain unlimited PIN levels. Thus, in this model cells can become polar or apolar depending on their initial auxin and PIN concentrations (Fig. 3B).
with their left neighbour, a circular auxin flux arises that results in a pattern in which final auxin levels are the same across all cells (Fig. 3C). Together, this demonstrates that the model displays polarity-driven self-organisation.

For a small range of parameters, linear feedback also permits two membrane segment equilibria (Fig. 3A), four single-cell equilibria (Fig. 3B) and self-organised polarity-driven patterning (Fig. 3C), as for quadratic feedback. However, for most parameter values, the upper unstable membrane segment equilibrium, and the domain of unlimited PIN levels that it demarcates, do not occur and instead only a single stable equilibrium will be produced (Fig. 3D). Under these conditions, the single-cell level model will have a single stable equilibrium, resulting in apolar cells (Fig. 3E). In the absence of cell polarity, the tissue level pattern now strongly depends on the persistence of sources and sinks and is not able to sustain itself after a temporal perturbation. Hence, the model behaves in a gradient-driven rather than a self-organising manner (Fig. 3F).

We conclude that the Stoma et al. (Stoma et al., 2008) model has the ability to self-organise for both linear and superlinear feedback, although self-organisation occurs in a considerably broader parameter range for the superlinear than the linear feedback. Note that it was analytically shown by Mitchison (Mitchison, 1980) that the formation of distinct veins rather than laminar flows in 2D tissue poses the more stringent requirement for superlinear feedback. This additional requirement for 2D symmetry breaking to generate distinct veins cannot be recovered using our 1D framework.

Altering the model to include a limiting PIN pool

A much discussed aspect of most flux-based models is the low auxin concentrations that they produce in veins (Kramer, 2008; Rolland-Lagan and Prusinkiewicz, 2005), which disagrees with the experimental data (Scarpella et al., 2006). The (non-physiological) cell level assumption of a PIN pool that is so large that membrane segments do not compete for it (termed ‘unlimited’) can be the cause of low auxin levels in the veins, as it allows for an unlimited increase in membrane PIN levels in response to auxin flux, and hence an unlimited efflux of auxin out of these flux channels.

By incorporating a limiting PIN pool (see Box 4) into the Stoma et al. (Stoma et al., 2008) model, maximum membrane PIN levels become limited by the total amount of PINs that a cell contains. Note that a similar effect can be obtained by assuming a saturating feedback function, which limits the level of PINs that auxin can induce on a membrane segment. At the membrane segment level, the addition of a finite PIN pool causes the model to become truly bistable (Fig. 3G), with a stable high equilibrium as opposed to a limited PIN pool behaviour for the superlinear than the linear feedback. Note that it was analytically shown by Mitchison (Mitchison, 1980) that the formation of distinct veins rather than laminar flows in 2D tissue poses the more stringent requirement for superlinear feedback. This additional requirement for 2D symmetry breaking to generate distinct veins cannot be recovered using our 1D framework.

At the tissue level, the model with a limiting PIN pool behaves similarly to the model without a limiting PIN pool (Fig. 3C), the difference being that polar cells now have a low PIN concentration on one membrane segment but a finite, high PIN concentration on the other. Thus, adding a finite PIN pool eliminates the unlimited growth of PIN concentrations and thus limits the maximal flux out of the cells. This is sufficient to allow auxin concentrations to build up in veins, as was demonstrated by Feugier and Iwasa (Feugier and Iwasa, 2006).

**Box 4. Mathematical analysis of the flux-based model of Stoma et al. (Stoma et al., 2008)**

**Membrane segment level**

The auxin equation is given by Eq. 1. The PIN equation for this model depends on the feedback of auxin flux on membrane PIN levels. Flux ($F$) consists of influx and passive and active efflux over the membrane segment:

$$ F = i_{aux} + i_{pass} - i_A. $$

Note that we use the same terminology as in the auxin equation (Eq. 1). Since in the cell to which the membrane segment belongs auxin is constant, it is incorporated in $i_{aux}$ and $i_{pass}$. Feedback of the auxin flux on membrane segment PIN levels is modelled through an increased recycling rate of PINs ($k_{on}$). Two alternative feedback functions have been proposed by the authors: the linear function $k_{on}=k_{on1}+k_{on2} F$ and the quadratic function $k_{on}=k_{on0}+k_{on2} P^2$. In both cases, the feedback only takes place when the flux is larger than 0, i.e. when there is net efflux. See Table 2 for parameter definitions and default values. Incorporating the linear feedback function in the $dP/dt$ equation (Eq. 3) gives the PIN equilibrium line:

$$ P = \begin{cases} k_{on} + k_{on1} (i_{aux} - e_A) & \text{if } F > 0, \\ k_{on1} - k_{on2} i_{aux} & \text{if } F \leq 0, \end{cases} $$

which gives the dashed line in Fig. 3A or 3D. This line can intersect twice with the auxin equilibrium line given by Eq. 2 (but see Fig. 3D). The lower equilibrium is stable and the upper is unstable. Above this equilibrium, unlimited increase of PINs and auxin takes place. The PIN equilibrium line with quadratic feedback gives similar results (dotted line in Fig. 3A).

**Single-cell level**

At the single-cell level, the authors assume the availability of PINs to be non-limiting. Hence, $P_0$ and $P_1$ can be described by Eq. 4b and, assuming that auxin concentrations are in dynamic equilibrium (Eq. 1), PIN equilibrium lines are given by:

$$ P_i = \frac{k_{on} (i_{aux} - e_A) + k_{pass} (e_A + d))}{k_{off} (e_A + d) - k_{on2} i_{pass} d} \text{ with } i = 0, 1. $$

Eq. 9 indeed shows that $P_0$ and $P_1$ are independent of each other. The equilibrium lines are exactly horizontal and vertical (Fig. 3B, solid lines for $P_0$ and dashed lines for $P_1$), intersecting once in a stable apolar equilibrium and three times in unstable equilibria. Above and to the right of these unstable equilibria, unlimited growth of either $P_0$ or $P_1$ or both takes place.

**Adding a limiting PIN pool**

We include a limiting PIN pool by incorporating the feedback functions into Eq. 5 rather than into Eq. 4b. The resulting equilibrium lines have similar shapes for linear and quadratic feedback (Fig. 3G). Owing to the limiting PIN pool, the top part of the PIN equilibrium line curves to the right, allowing for two stable equilibria at the membrane segment level. At the single-cell level, the curves now obtain a complicated shape that allows them to intersect in three stable equilibria: two polar and an apolar rest state (Fig. 3H).
Example 2: concentration-based model
Smith et al. (Smith et al., 2006) formulated a model for phyllotaxis in the growing shoot apical meristem. Membrane PIN levels were assumed to depend positively on the auxin levels in neighbouring cells. The authors used a superlinear saturating feedback function (see Glossary, Box 1), meaning that at first PIN levels increase more than linearly with auxin concentrations until maximum membrane PIN levels are reached. In addition, a limited PIN pool is assumed, causing a competition for PINs between the different membrane segments of a cell. Finally, auxin pumping by the PINs is assumed to saturate with auxin levels in the neighbouring cell. The authors show that, for an initially homogenous ring of cells, evenly spaced peaks of auxin arise after the application of small, transient perturbations, with PINs pointing toward the neighbouring cell with the highest auxin level. In a 2D epidermal tissue layer, the model also produced distinct peaks; however, this occurred at variable distances from one another. However, after incorporating differences in auxin handling between the cell types of different meristem domains, the model can produce stable phyllotaxis patterns. The authors furthermore showed that distinct phyllotaxis patterns can be reproduced by the model depending on parameter values such as meristem size and the transport and diffusion of auxin.

Our analysis of the model by Smith et al. (Smith et al., 2006) is described in Box 5. We find that, at the membrane segment level, the model indeed allows for bistable behaviour (Fig. 4A). Combined with the limiting PIN pool that causes competition for PINs between membrane segments, this results in a cell model with two stable equilibria in which either $P_0$ or $P_1$ is low or vice versa (Fig. 4B). Hence, cells in this model display polar behaviour. Our analysis of the model by Stoma et al. (Stoma et al., 2008) after incorporating a finite PIN pool, in which a stable apolar rest equilibrium was also present (Fig. 3H).
Box 5. Mathematical analysis of the concentration-based model of Smith et al. (Smith et al., 2006)

Membrane segment level
To model auxin dynamics, Smith et al. (Smith et al., 2006) take into account production, decay, influx and efflux processes. Auxin production is assumed to be limited by the amount of auxin already present in the cell. Auxin efflux depends on a saturating manner on auxin content in the neighbouring cell, thus both influx and efflux are dependent on auxin in the neighbouring cell (A) and auxin in the cell to which the membrane segment of interest belongs (A_i):

\[
\frac{dA}{dt} = \frac{p}{1 + kA_i} - \frac{dA + e_{pin}A_i}{h^2_{pin} + A_i^2} - \frac{e_{pin}A_i}{h^2_{pin} + A_i^2} + e_{pas} A_i - \frac{e_{pas}A_i^2}{h^2_{pin} + A_i^2} \quad \text{with } i = 0, 1,
\]

in which \( k \) is the factor by which auxin inhibits its own production, \( e_{pin} \) is the passive efflux out of the neighbouring cell and \( e_{pas} \) is the PIN-mediated efflux, which is half maximum when \( A = h_{pin} \). Note that both \( A \) and the PINs on the membrane of a neighbouring cell are still assumed to be constant. The resulting auxin equilibrium line is given by:

\[
P_i = \frac{h_{pin}^2 + A_i^2}{e_{pin} A_i^2} \left( A_i + e_{pin} A_i^2 \right) - \frac{p}{1 + kA_i} - e_{pas} \quad \text{with } i = 0, 1.
\]

Smith et al. (Smith et al., 2006) assume the PIN pool to be limited, all PINs to reside on membranes and PIN dynamics to always be in equilibrium. Thus, they determine the PIN equilibrium line directly:

\[
P = \frac{P_{tot} b^A}{b^A + b^A}.
\]

where \( b \) is a base parameter that the authors set to 2 or 3, and the sum is taken over the auxin concentrations in all neighbouring cells (\( n \)). \( P_{tot} \) contains explicit (auxin-induced) production and decay of PINs. Given that we maximally consider two auxin concentrations in neighbouring cells, we can replace Eq. 12 with:

\[
P_i = \frac{P_{tot} b^A}{b^A + b^A} \quad \text{with } i = 0, 1.
\]

Note that, at the membrane segment level, we assume \( A_i \) to be constant. There are a total of three intersection points between the PIN and auxin equilibrium lines (dashed and solid lines in Fig. 4A), of which the outer two equilibria, corresponding to low and high PIN levels, are stable.

The PIN equilibrium line shifts with the value of \( b^A \). It moves to the right when \( A_i \) is high and to the left when \( A_i \) is low (Fig. 4D,E, respectively), potentially eliminating two equilibria and thus the potential for bistability. Thus, whether cells can polarise depends on their local auxin context.

Single cell
For the cell level, the PIN equilibrium line is given by Eq. 13, in which now both \( A_i \) and \( A_j \) are variable. We substitute these in the auxin equations (Eq. 10) and solve to obtain the auxin equilibrium lines in Fig. 4B. These intersect three times in two stable polar equilibria separated by an unstable one. Since the PINs are assumed to be in equilibrium with the auxin levels (Eq. 13), they can be directly deduced from the auxin levels in the equilibria.

Removing the limiting PIN pool
To study how the model behaviour changes if there is no competition for a limited PIN pool, we have to reverse-engineer Eq. 12. The PIN equation is given by Eq. 5 and there are two possible ways to create the sigmoidal PIN equilibrium line through feedback of auxin on \( k_{on} \). In a non-saturating (\( k_{on}=k_{on}+k_{off}b^A \)) or a saturating (\( k_{on}=k_{on}+k_{off}b^A(b^A+b^A+b^A) \)) function. Substituting these feedbacks into the PIN equation without a limiting PIN pool (Eq. 3) results in the PIN equilibrium lines:

\[
P = \frac{1}{k_{on}} \left( k_{on} + k_{off} b^A \right),
\]

\[
P = \frac{1}{k_{on}} \left( k_{on} + k_{off} b^A \right).
\]

Both of these lines can intersect more than once with the auxin equilibrium line. In the case of Eq. 14a there are two equilibria (Fig. 4F). The lower one is stable, and the upper, which is unstable, separates a region of unlimited increase in PINs. In the case of Eq. 14b, the system has two stable equilibria separated by an unstable one (Fig. 4H). At the cell level this results in one stable equilibrium and three regions of increase of either \( P_0 \), \( P_1 \), or both for Eq. 14a (Fig. 4G) or four stable equilibria, two polar and two apolar, for Eq. 14b (Fig. 4I).

In a file of cells, we find that PINs are oriented up-the-gradient (Fig. 4C), which is consistent with the premise of the model whereby PINs are localised on the membrane segment apposing the neighbouring cell with highest auxin levels. Polarised cells display a low PIN concentration on the membrane segment facing the sink and a high PIN concentration on the membrane facing the source, and these levels are independent of the strength of the source and sink. When a ring of cells is simulated, a transient perturbation in one of the cells spreads throughout the entire cell file, producing a persistent polarisation of all cells. The cell that received the perturbation, in the form of a transient increase in auxin, develops into an auxin minimum at the opposing side of the cell file, from which the polarised cells point away. Thus, in contrast to what happens in the flux-based model, the polarisation of cells is accompanied here by a patterned rather than homogeneous auxin distribution (Fig. 4C). Our analysis thus confirms that pattern formation in the Smith et al. (Smith et al., 2006) model is self-organising, and that this self-organisation is polarity-driven, in agreement with the model behaviour reported in the original paper.

The feedback function used by Smith et al. (Smith et al., 2006) results in a PIN equilibrium line (see Glossary, Box 1) that shifts as a function of the difference in auxin concentration between two neighbouring cells (see Fig. 4D,E). As a result, one of the two membrane segment equilibria may disappear, abolishing membrane bistability and cell polarity. Therefore, not all cells within a tissue necessarily experience the auxin concentration differences that are necessary to polarise. This might explain why the authors observed strongly polarised cells close to maxima, but weaker or no polarisation at a greater distance from the maxima.

Removing the limiting PIN pool from the model
Next we examine whether the model behaviour changes if we assume that there is no competition between membrane segments for a limiting PIN pool (see Box 5 for details). Obviously, this is not increasing the physiological realism of the model. However, performing this analysis allows us to compare the flux-based and concentration-based models both under conditions with and without a finite PIN pool. On the membrane segment level, we find that there is still bistability, with either two stable and one unstable equilibrium as in Fig. 4A (Fig. 4H, for saturating feedback) or with one stable and one unstable equilibrium (Fig. 4F, for non-saturating feedback). On the cell level, we find four alternative equilibria for non-saturating feedback (Fig. 4G), which result from combining the two
equilibria for $P_0$ with the two equilibria for $P_1$. The single stable equilibrium corresponds to the apolar cell state in which both membrane segments have a low PIN level. The unstable equilibria separate this apolar rest state from the alternative states in which unlimited growth of PINs on either one or both of the membrane segments occurs. If, instead, the feedback is saturating, there are a total of nine equilibria at the single-cell level (Fig. 4H), which result from combining the three equilibria at the membrane segment level. In this case there are four stable equilibria: one apolar, two alternative polar equilibria, and one bipolar. Thus, despite the removal of the PIN pool the model retains its potential to produce polar cells. However, it acquires a bipolar state, with high PINs on both membranes or a region in which both membranes increase their PIN level unlimitedly, and an apolar rest state in which both membranes have low PIN levels. Consequently, the robustness of polarity-driven self-organised patterning is decreased. Thus, the tissue level behaviour is similar upon removal of the PIN pool, but whether a cell is polar or not now depends more strongly on its context.

**Evaluation of other PAT models**

Having analysed these two example models in detail, we now briefly discuss our analysis of other PAT models found in the literature. The full analysis can be found in supplementary material Appendix S1, and an overview of all model assumptions and the generated behaviour is given in Table 2.

**Flux-based models**

In addition to the model by Stoma et al. (Stoma et al., 2008) that we used as an example, we analysed a variety of published flux-based models (Mitchison, 1980; Mitchison, 1981; Feugier and Iwasa, 2006; Feugier et al., 2005; Fujita and Mochizuki, 2006; Alim and Frey, 2010). Summarising, we find that all flux-based models, independent of the shape of the feedback function, auxin pumping dynamics or whether a limiting PIN pool is assumed, are capable of generating polarity-driven self-organised auxin and PIN patterns. Furthermore, all these models contain a stable rest state equilibrium at the single-cell level, in which both membranes contain low PIN levels, that is separated from the two alternative
polar states by unstable equilibria. Consequently, cell polarisation in these models requires perturbations that push the system beyond these unstable equilibria into the realm of the polar equilibria.

We find that in those models (Feuger and Iwasa, 2006; Feuger et al., 2005; Fujita and Mochizuki, 2006; Alim and Frey, 2010), in which a finite PIN pool is assumed, no bipolar equilibrium with high PIN levels at both membranes is present. By contrast, in the early models by Mitchison (Mitchison, 1980; Mitchison, 1981) and in the model by Stoma et al. (Stoma et al., 2008) (which is based on these early models) no limiting PIN pool is incorporated and a bipolar state does occur, resulting in more possibilities for nonpolar cells. Similar to the model by Stoma et al. (Stoma et al., 2008), the other flux-based models show with-the-flux behaviour, but can generate up-the-gradient auxin transport depending on the strength of localised sinks. In all cases, we find that a superlinear dependence of membrane PIN levels on auxin flux increases the parameter range for membrane bistability and hence the robustness of self-organised patterning.

Concentration-based models
Besides the model by Smith et al. (Smith et al., 2006), we applied our analysis to several further published concentration-based models (Jönsson et al., 2006; Merks et al., 2007; Newell et al., 2007; Sahlin et al., 2009). For the model by Newell et al. (Newell et al., 2007), which combines auxin concentration feedback on PINs and mechanical effects, we have not considered the mechanical effects. Concentration-based models do not automatically allow for self-organised patterning. For polarity-driven self-organised patterning to arise, either a non-linear feedback function (Smith et al., 2006; Jönsson et al., 2006) and/or non-linearity in auxin pumping dynamics (Jönsson et al., 2006; Merks et al., 2007; Sahlin et al., 2009) is necessary. This difference between concentration-based and flux-based models is due to the inherent presence of non-linearity in flux-based models. First, in most flux models, feedback only occurs for net efflux, with PIN levels abruptly switching to the minimum concentration for negative efflux. Hence, even a seemingly linear function is transformed into a non-linear function. Second, flux is essentially the product of the PIN level and auxin concentration. Thus, even if flux feeds back linearly on PINs, this results in a direct positive feedback of PINs on PINs, effectively making the feedback non-linear. Increased non-linearity in feedback or auxin pumping, saturating feedback and a finite PIN pool contribute to the robustness of polarity-driven self-organised patterning, as was the case in flux-based models. Introduction of a limiting PIN pool in concentration-based models abolishes both the bipolar and the apolar rest equilibrium, whereas in flux-based models only the bipolar equilibrium disappears. Consequently, concentration-based models incorporating a finite PIN pool always produce polar cells without requiring any substantial perturbation, as no stable apolar cell state is present.

In addition to polarity-driven self-organised patterning, we find an additional self-organisation mechanism in concentration-based models that occurs if a finite PIN pool is combined with a linear feedback function and linear auxin pumping dynamics (Newell et al., 2007; Jönsson et al., 2006). Owing to the absence of non-linearity, no polarity-driven self-organisation can arise. However, on the tissue level, auxin and PIN polarity patterns look similar, as for the polarity-driven self-organisation discussed for Smith et al. (Smith et al., 2006). The principal difference is that, in isolation, single membrane segments are not bistable and single cells are not polar, and hence polarity arises only at the tissue level. This mechanism of self-organisation resembles Turing-type patterning. The positive feedback from cellular auxin levels on PIN levels in the membranes of neighbouring cells amplifies local auxin maxima, thus functioning as short-range activation, while the resulting depletion of auxin from surrounding tissue prevents the formation of nearby maxima, thus serving as long-range inhibition. For certain parameter conditions, and combined with a finite PIN pool that causes competition for PINs between membranes and hence promotes cellular polarity, these effects destabilise the single stable equilibrium of the system and lead to the formation of auxin maxima and PIN polarity (see our bifurcation analysis in supplementary material Appendix S1). The detailed analysis by Sahlin et al. (Sahlin et al., 2009) suggests that this mechanism might, in addition to the isolated maxima needed for phyllotactic patterns, also generate stripe-like patterns.

Finally, in contrast to flux-based models, concentration-based models are capable only of generating up-the-gradient patterning.

Mechanistic models
An alternative hypothesis for PIN polarisation was introduced by Heisler et al. (Heisler et al., 2010). The model was inspired by the observation that PIN polarity is correlated with the alignment of cortical microtubules. This led the authors to propose that PINs localise to membrane segments adjoining cell walls that experience the most stress, which in turn is the result of local cell expansion due to auxin. The model can thus be seen as a more mechanistic version of concentration-based models, with wall stress being the readout of auxin content in the neighbouring cells. The combination of non-linear auxin pumping dynamics and a finite PIN pool causes the model to behave in a self-organised polarity-driven manner that generates up-the-gradient patterning, similar to concentration-based models.

Recently, Wabnik et al. (Wabnik et al., 2011) proposed yet another alternative hypothesis for the feedback of auxin on membrane PIN levels. In their model, extracellular auxin binds to free receptors in the apoplast. The resulting auxin-receptor complexes inhibit endocytosis of PIN proteins on the nearest membrane segment. Furthermore, the complexes limit the diffusion of auxin and receptor, producing both an intra-apoplast auxin gradient and competition for receptors between segments on either side of the cell wall. We extended our framework with cell wall compartments and receptor variables in order to study this model (see supplementary material Appendix S1). Our analysis shows that the model allows for membrane bistability and cell polarity and produces with-the-flux tissue polarisation patterns under the additional requirement that auxin has a low diffusion speed in the cell wall, hence allowing for wall gradients.

Models for fountain-like patterns
For the model by Stoma et al. (Stoma et al., 2008), we discussed how the authors used the flux-based mechanism to generate both up-the-gradient maximum formation in the epidermis and down-the-gradient with-the-flux vein formation in subepidermal tissues. In order to achieve this, distinct model settings were required for the different tissue layers. Similarly, Merks et al. (Merks et al., 2007) proposed a concentration-based model to simulate these two processes in combination. In addition to the assumption that cells position their PINs toward the neighbouring cell with the highest auxin concentration, auxin is assumed to induce PIN expression (Vieten et al., 2005), which destabilises the position of the auxin maximum. As a result, the initially epidermal maximum propagates.
Table 2. Evaluation of different models according to assumptions and behaviour on all levels of organisation

<table>
<thead>
<tr>
<th>Reference</th>
<th>Mechanism</th>
<th>Feedback function</th>
<th>Biological interpretation</th>
<th>Auxin transport</th>
<th>Biological interpretation</th>
<th>PIN dynamics</th>
<th>PINs cycles</th>
<th>PINs organised</th>
<th>Polar cells?</th>
<th>Maximal organisation</th>
<th>Up or down gradient</th>
</tr>
</thead>
<tbody>
<tr>
<td>Stoma et al., 2008</td>
<td>Flux-based</td>
<td>$k_{on} = F$</td>
<td>Exocytosis increases linearly with flux</td>
<td>Linear: $T = PA$</td>
<td>$\frac{dP}{dT} = k_{on} - k_{off}P$</td>
<td>PINs cycle to and from the membrane; exocytosis does not lower the availability of cytosolic PINs</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Up or down gradient</td>
<td></td>
</tr>
<tr>
<td>Smith et al., 2006</td>
<td>[I]-based</td>
<td>$P = \frac{b^n}{\sum_n}$</td>
<td>[PIN] increases faster with higher [auxin] in neighbour, and saturates for high [auxin]</td>
<td>Saturated: $T = P \frac{A^2}{h^2 + A^2}$</td>
<td>PINs is limiting; transport saturates for high [auxin] in neighbouring cell</td>
<td>PIN pool is limiting, all PINs are on the membrane, and cycling is in equilibrium; auxin induces PIN expression</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Up</td>
<td></td>
</tr>
<tr>
<td>Mitchison, 1980</td>
<td>Flux-based</td>
<td>$D = \frac{F}{h^2 + F}$</td>
<td>Permeability increases more with flux when flux becomes higher</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Mitchison, 1981</td>
<td>Flux-based</td>
<td>$D = \frac{F}{h^2 + F}$</td>
<td>Permeability increases more with flux when flux becomes higher</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Feugier et al., 2005</td>
<td>Flux-based</td>
<td>Various</td>
<td>Various</td>
<td>Various</td>
<td>Various</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Feugier and Iwasa, 2006</td>
<td>Flux-based</td>
<td>$k_{on} = F$</td>
<td>Exocytosis increases linearly with flux when flux becomes higher</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Fujita and Mochizuki, 2006</td>
<td>Flux-based</td>
<td>$P = \frac{e^A}{1 + e^A}$</td>
<td>[PIN] increases linearly with [auxin] in neighbour and saturates for high [auxin]</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Aim and Frey, 2010</td>
<td>Flux-based</td>
<td>$k_{on} = F$</td>
<td>Exocytosis increases linearly with flux when flux becomes higher</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Down</td>
<td></td>
</tr>
<tr>
<td>Jonsson et al., 2006</td>
<td>[I]-based</td>
<td>$P = \frac{A}{\sum_n A_n}$</td>
<td>[PIN] increases linearly with [auxin] in neighbour and saturates for high [auxin]</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>No</td>
<td>No</td>
<td>Self-organised: other (see text)</td>
<td>Up</td>
<td></td>
</tr>
<tr>
<td>Merks et al., 2007</td>
<td>[I]-based</td>
<td>$k_{on} = \frac{A}{h + A}$</td>
<td>Exocytosis increases linearly with [auxin] in neighbour and saturates for high [auxin]</td>
<td>Saturated: $T = P \frac{A}{h + A}$</td>
<td>PIN is limiting; transport saturates for high [auxin]</td>
<td>PINs are produced, decayed and cycle to and from the membrane; auxin induces PIN expression</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Up</td>
<td></td>
</tr>
<tr>
<td>Newell et al., 2007</td>
<td>[I]-based</td>
<td>$k_{on} = A$</td>
<td>Exocytosis increases linearly with [auxin]</td>
<td>Linear: $T = PA$</td>
<td>PIN is limiting; transport saturates for high [auxin]</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting</td>
<td>No</td>
<td>No</td>
<td>Self-organised: other (see text)</td>
<td>Up</td>
<td></td>
</tr>
<tr>
<td>Sahlin et al., 2009</td>
<td>[I]-based</td>
<td>Various</td>
<td>Various</td>
<td>Saturated: $T = P \frac{A}{h + A}$</td>
<td>PIN is limiting; transport saturates for high [auxin]</td>
<td>PINs cycle cycled from the membrane, PIN pool is limiting (some tests performed with an unlimited PIN pool)</td>
<td>Yes</td>
<td>Yes</td>
<td>Self-organised: polarity-driven</td>
<td>Up</td>
<td></td>
</tr>
</tbody>
</table>
subepidermally through the tissue and leaves basally pointing PINs in its wake. Thus, whereas the model by Stoma et al. (Stoma et al., 2008) required different assumptions for different tissue layers to combine auxin maximum and vein formation, the model by Merks et al. (Merks et al., 2007) is only capable of generating the two patterns sequentially, with one replacing the other.

In another attempt to explain fountain-like pattern formation, Bayer et al. (Bayer et al., 2009) developed a model that explicitly combines concentration-based and flux-based feedback, using an auxin level threshold to decide which type of feedback should be applied locally. The model generates auxin maxima in the epidermis by invoking up-the-gradient feedback and then switches to with-the-flux feedback that positions PINs in underlying tissues away from the maximum. In the up-the-gradient regime, the authors chose the same feedback function as in their previous concentration-based model (Smith et al., 2006). Hence, the membrane segments are bistable and single cells are polar, allowing polarity-driven self-organised pattern formation. In the with-the-flux regime, the feedback function is superlinear. As shown by Feugier et al. (Feugier et al., 2005) and following from our analysis, these conditions also allow for polar cells and polarity-driven self-organising vein formation with high auxin concentrations in the veins. However, given that two distinct feedback mechanisms need to be invoked, the overall pattern is not fully self-organised.

**Extrapolating our analysis to 2D and 3D tissue and in vivo plant patterns**

Our analysis is restricted to the self-organising potential of published models at the membrane segment, single-cell and 1D tissue level and the type (orientation) of 1D patterns that they generate. The reason for this restriction is mainly the feasibility of a rigid analytical approach. Obviously, this limits the degree to which we can predict the 2D and 3D tissue patterns generated by the models and how far these model patterns correspond to in vivo patterns. For example, based on our analysis we cannot predict the precise spacing and arrangement between phyllotactic maxima or leaf veins. Indeed, Mitchison (Mitchison, 1980) showed analytically that to obtain an additional symmetry breaking in 2D tissue and generate distinct veins rather than laminar flow patterns, superlinear feedback of auxin flux on PIN polarisation is required. Based on our 1D analysis, we can only conclude that both linear and superlinear feedback generate self-organised-with-the-flux oriented polarisation patterns, but cannot determine this additional demand for distinct veins. In other words, we can determine the necessary requirements for self-organised patterning, but we cannot determine requirements for a specific phyllotactic or venation pattern. However, we are able to extrapolate our 1D analysis to higher dimensional tissue patterns to a qualitative extent, focusing on the orientation and self-organised nature of the patterning.

Let us consider the fountain and reverse-fountain types of flux pattern observed in the shoot primordia and root tip (Fig. 1C,D). Two main transport directions can be distinguished in both organs: toward (labelled 2) and away (labelled 1) from a maximum. In the root tip (Fig. 1D), auxin flux is directed toward the root tip maximum in the inner tissue layers and away from the maximum in the outer layer(s). The pattern in a shoot primordium is reversed (Fig. 1C). Our 1D analysis reveals how, in each model, PINs can polarise with respect to a maximum. We found that concentration-based models always direct PINs toward an auxin maximum, whereas flux-based models generally locate PINs away from the maximum but can be made to display an opposite orientation under particular parameter conditions. Hence, concentration-based models can explain the arrows labelled 2 in Fig. 1C,D, which point towards the auxin maximum, but cannot explain those labelled 1 in an autonomous manner. However, it has been shown that expression of the gene *PINOID* can switch the polar orientation of PINs from a basal to an apical orientation (Friml et al., 2004) by influencing intracellular PIN trafficking (Michniewicz et al., 2007). Thus, in the root, PINOID expression in the outer tissue layer(s) could account for the switch to away-from-maximum localisation in these layers. However, this would not explain the away-from-maximum arrows in the inner tissues of the shoot primordium. Flux-based models can in theory account for both arrows. However, this requires that the auxin maximum serves as a sink (that attracts auxin) for the inner layers and as a source for the outer layers in the root tip (with the reverse for the shoot primordium) [see model by Stoma et al. (Stoma et al., 2008)].

Thus, based on our analysis we conclude that, despite their ability to generate self-organised patterns, neither concentration-

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### Table 2. Continued

| Bayer et al., 2009 | Flux-based and [l]-based | $P = \frac{b}{\sum b_i}$ | $[\text{PIN}]$ increases faster with higher flux/[auxin] in neighbour and saturates for high flux/[auxin] | Saturated: | $T = P = \frac{A^2}{h^2 + A^2}$ | $[\text{PIN}]$ is limiting: transport saturates for high [auxin] | $P = \frac{P_k b_i}{\sum b_i}$ | PINs produced and decayed in $P$, PINs are produced, decayed; cycling is in dynamic equilibrium and all PINs are on the membrane; auxin induces PIN expression | Yes | Yes | Self-organised: polarity-driven | Up and down |

| Hesler and Jönsson, 2006 | [l]-based | $P = \frac{A}{\sum b_i}$ | Endocytosis increases linearly with [auxin] in neighbour and saturates for high [auxin] | Saturated: | $T = P = \frac{A}{h + A}$ | $[\text{PIN}]$ is limiting: transport saturates for high [auxin] | $P = \frac{P_k b_i}{\sum b_i}$ | PIN pool is limiting, all PINs are on the membrane, and cycling is in dynamic equilibrium | Yes | Yes | Self-organised: polarity-driven | Up |

| Wabnik et al., 2011 | Cell wall [l]-based | $k_{\text{off}} = \frac{1}{h} A$ | Endocytosis decreases linearly with [auxin] in cell wall and saturates for high [auxin] | Saturated: | $T = P = \frac{A}{h + A}$ | $[\text{PIN}]$ is limiting: transport saturates for high [auxin] | $\frac{dP}{dt} = -k_{\text{off}} P - k_{\text{off}} P$ | PINs produced, decayed, and cycle to and from the membrane; auxin induces PIN expression | Yes | Yes | Self-organised: polarity-driven | Up |
based nor flux-based models can explain fountain-type patterns in a fully autonomous self-organised manner, but require additional tissue- or location-specific assumptions. This implies either that auxin patterning in plants is not fully self-organised and that tissue-specific responses are required, or that the currently proposed feedback mechanisms need to be revised to fully explain self-organised auxin patterning.

Conclusions
Auxin patterning plays a central role in robust, yet developmentally and environmentally flexible, plant development. The long-standing idea that auxin influences its own transport, thus potentially allowing for self-organised patterning, has inspired numerous modelling studies. Apart from the major and easily understood distinction assumed in the feedback mechanism – whether auxin flux across the membrane or auxin levels in neighbouring cells impacts membrane PIN levels – models differ in their mathematical descriptions of the details of this feedback, as well as of auxin transport dynamics and PIN cycling dynamics. The relevance of these variations for model output is less obvious. In addition, little attention has been devoted to whether the generated model patterns arise in an autonomous self-organised manner or are strongly dependent on persistent prepatterns or additional assumptions. However, largely self-organised patterning is essential for robust, repeated and flexible patterning during plant development. Here we developed a generalised analytical and simulation framework to compare most of the currently published PAT models, translating the wide range of mathematical functions employed into a limited set of underlying biological assumptions, and reporting the model behaviour that this results in on the membrane segment, cell and tissue levels. We focused on the directionality of PIN polarisation and the extent of self-organised patterning that the models generate.

To summarise our analysis, flux-based models are somewhat more flexible in the direction of polarisation that they can generate, and do not necessarily require non-linearity or a limiting PIN pool to generate self-organised patterning. By contrast, concentration-based models allow for two different modes of self-organised patterning, and require less perturbation for cell polarisation to occur in the case of a limiting PIN pool. Still, both model types require largely similar biological assumptions to generate patterns in a robust self-organising manner, and both display only a single strongly preferred direction of polarisation. However, in plants, opposite polarities often co-occur. In the shoot and lateral root primordia, leaves, the primary root and the developing embryo, (reverse) fountain-like patterns are found, with neighbouring cell files displaying opposite PIN polarisations with respect to the same tissue gradient. Based on our analysis, we therefore hypothesise that neither concentration-based nor flux-based models are currently capable of explaining bidirectional fountain-type patterns in a fully self-organised manner.

Future directions
In this exploration of published PAT models, we have demonstrated that all current models explain in planta auxin and PIN patterning to the same limited extent and for similar conditions. Consequently, based on our current knowledge it remains unclear which feedback mechanism is at work in real plants – concentration-based, feedback-based, a combination of the two, or an alternative mechanism. Further, it is not clear how these models can explain the robust and repeated formation of fountain-type patterns in primordia and apical and basal meristems. In which directions should future modelling and experimental research proceed in order to answer these questions?

Let us first focus on the modelling. Our analysis shows that a positive auxin-PIN feedback, combined with a non-linearity in either feedback or auxin pumping or a limiting PIN pool, suffices for generating robust self-organised patterning. Thus, we would argue that further extending the number of flux-based and concentration-based models (by varying the type of mathematical functions used for describing the biological processes and by varying how the required non-linearity is incorporated) will not advance the current situation, but will simply provide more of the same results. However, this knowledge does allow us to easily generate alternative feedback models – for example, a feedback of local, intracellular auxin concentration on membrane PIN levels as suggested by Kramer (Kramer, 2009) – that have self-organising capacity, and to study to what extent such alternative feedback models are capable of generating fountain-type patterns in a self-organised manner.

Second, we discussed how tissue-specific PINOID expression combined with a concentration-based feedback, or a differential source and sink behaviour of an auxin maximum for different tissue layers combined with a flux-based model (Stoma et al., 2008), may be capable of generating at least certain fountain-type patterns. However, as these tissue-specific properties are superimposed, this patterning is not fully self-organised. Although only partly successful in the model by Merks et al. (Merks et al., 2007), a possibly fruitful approach might be to incorporate gene expression regulation into PAT models in such a manner that the tissue-specific requirements for different flux orientations are automatically generated as part of the patterning process. Important candidate genes to consider are the PINs themselves, additional exporters such as the P-GLYCOPROTEINs (PGPs), but also auxin importers such as the AUX/LAX genes, and PIN cycling dynamics regulators such as PINOID (e.g. Swarup et al., 2000; Bandyopadhyay et al., 2007; Benjamins et al., 2001). Together, these two approaches will hopefully allow us to determine the type of feedback, non-linearities and gene regulation mechanisms that are theoretically capable of generating self-organised fountain patterns.

To establish the mechanisms at work in real plants a combined experimental and modelling approach is necessary. Thus, a third important direction for future research is to formulate more molecularly mechanistic PAT models that allow for explicit experimental verification. This applies to tissue properties, such as cell walls and subcellular compartments, that have been ignored in a subset of the current models, but most importantly to the mechanism by which auxin feeds back on PIN localisation. The field has recently begun to move in this direction with models (Heisler et al., 2010; Wabnik et al., 2010) in which cells measure wall stress or use the ABP1 receptor to sense apoplast auxin levels, rather than measuring auxin levels in neighbouring cells. Although most molecular mechanisms behind the interactions proposed by Heisler et al. (Heisler et al., 2010) have yet to be determined, a property such as wall stress is compliant to experimental manipulation. Similarly, a local requirement for the ABP1 receptor, which plays a central role in the model by Wabnik et al. (Wabnik et al., 2010), can be tested experimentally. In addition, this model requires and predicts a steep auxin gradient in the apoplast. Ideally, spatiotemporally refined methods for measuring extracellular auxin concentrations should be established to test the existence of this gradient. One could also propose more mechanistic models for flux-based feedback, for example by assuming that PINs measure
the auxin efflux that occurs through them, rather than membranes measuring the total net efflux across them, which would be more amenable to experimental tests.

To enable the construction and refinement of these more mechanistic models, we need to gain a much better understanding of PIN cycling, expression and degradation, and how these might depend on auxin levels or fluxes. For example, our analysis points to the importance of a limiting cellular PIN pool for the robustness of self-organised patterning and for realistic auxin levels in veins. However, we also find that, in models that take PIN production and decay into account, the seemingly reasonable assumption that PIN degradation is limited to non-membrane-bound PINs implicitly causes the PIN pool to be non-limiting (Merks et al., 2007; Wabnik et al., 2010). Thus, we need to experimentally verify whether membranes compete for a finite cellular PIN pool by performing detailed quantification of the amount of PIN proteins present in different cellular compartments under different conditions. In addition, for these conditions, experiments using photo-convertible tags should be performed to determine PIN cycling rates to and from the membrane.

The proposed combining of, and iteration between, increasingly mechanistic PAT models and experiments, together with an extension of PAT models to encompass alternative feedback mechanisms and regulated gene expression, will allow us to eventually pinpoint the inner workings of auxin patterning in plants.

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